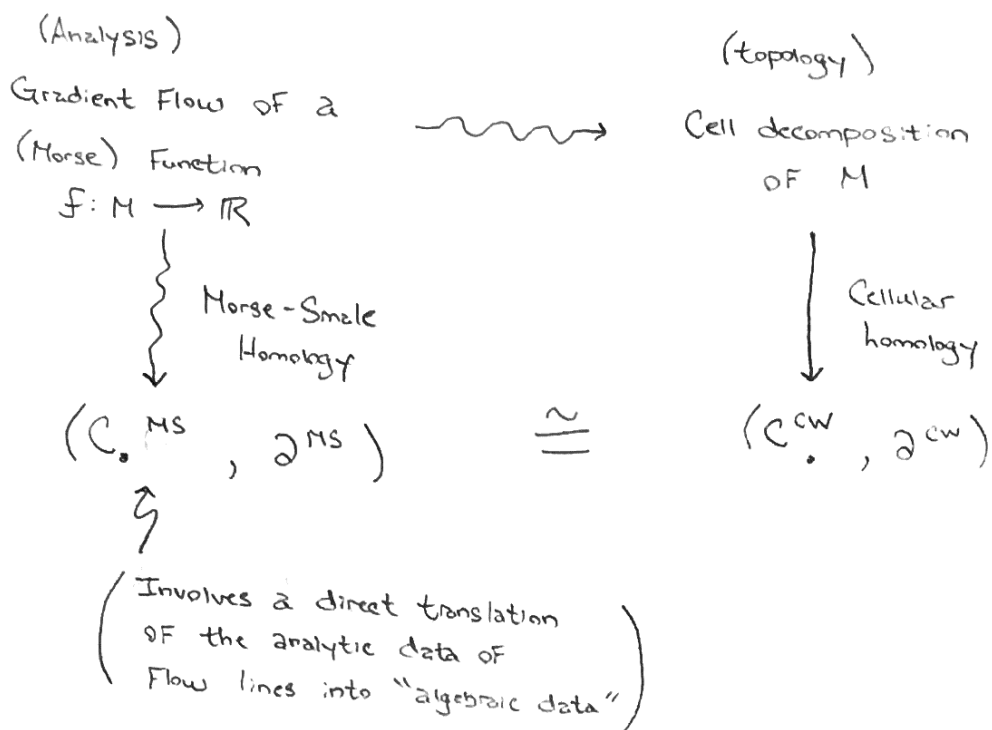


Recall. Morse TheoryWhat is the Morse-Smale Complex?Quick Defs

Setup: M smooth, cpt, n -dim \mathbb{R}
 $f: M \rightarrow \mathbb{R}$

Def

- $Crit(f) := \{c \in M \mid df_c = 0\}$
- f is Morse if $\forall c \in Crit(f)$; $Hess_c(f) = \left(\frac{\partial^2 f}{\partial x^i \partial x^j} \right)_{i,j=1}^n \Big|_{x=c}$ is non-degen.
- $c \in Crit(f)$; $Ind(c) := \#$ neg e-values of $Hess_c(f)$

Equip M w/ a metric g and study the negative gradient flow of f

$$\Phi_s^f: M \rightarrow M = \text{time } s \text{ Flow of } -\nabla f$$

Q: Local behaviour @ crit points?

Def: $p \in \text{Crit}(f)$

$$\mathcal{U}(p) := \left\{ m \in M : \lim_{s \rightarrow -\infty} \Psi_s(m) = p \right\} \leftarrow \text{(unstable)}$$

$$\mathcal{S}(p) := \left\{ m \in M : \lim_{s \rightarrow +\infty} \Psi_s(m) = p \right\} \leftarrow \text{(stable)}$$

Local picture around p : Use Morse Lemma.

\exists local coords $(x^i)_{i=1}^n$ s.t.

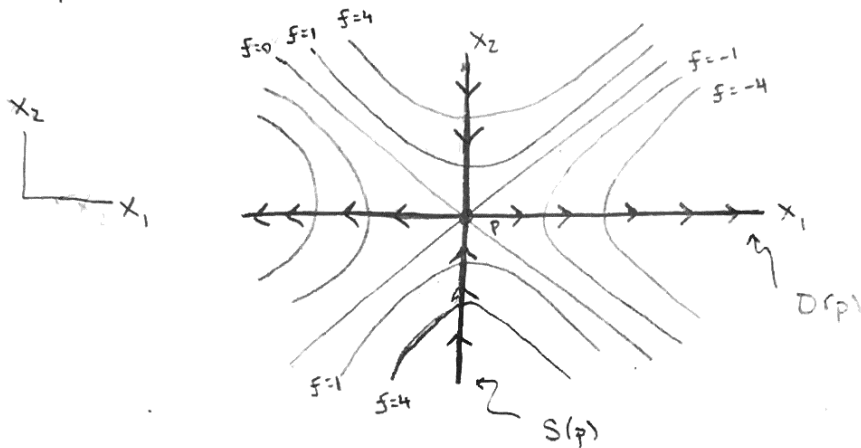
$$f = f(p) + -x_1^2 - \dots - x_{\text{ind}(p)}^2 + x_{\text{ind}(p)+1}^2 + \dots + x_n^2$$

\Rightarrow

$$\dim \mathcal{U}(p) = \text{ind}(p)$$

$$\dim \mathcal{S}(p) = n - \text{ind}(p)$$

Eg. $f = -x_1^2 + x_2^2$



Now define

$$\mathcal{M}(p, q) := \mathcal{U}(p) \cap \mathcal{S}(q) / \mathbb{R} = \left\{ \text{Flow lines } p \rightarrow q \right\} / \left(\begin{matrix} \text{Time} \\ \text{Translation} \end{matrix} \right)$$

Translation by $\Psi : \mathbb{R} \rightarrow \text{DIFF}(M)$

So (assuming Morse-Smale $\Rightarrow \mathcal{U}(p) \pitchfork \mathcal{S}(q)$)

$$\boxed{\dim \mathcal{M}(p, q) = \text{ind}(p) - \text{ind}(q) - 1}$$

Interested in when $\dim \mathcal{M}(p, q) = 0$! ($\text{ind}(q) = \text{ind}(p) - 1$)

DeF Morse-Smale Complex

$$C_i^{MS}(M, g, f) := \mathbb{Z} \langle \text{Crit}_i(f) \rangle \leftarrow \text{index } i \text{ critical points}$$

$$\partial^{MS} p = \sum_{q \in \text{Crit}_{i-1}} \# \mathcal{M}(p, q) \cdot q$$

Signed # (choose orientations for unstable manifolds)

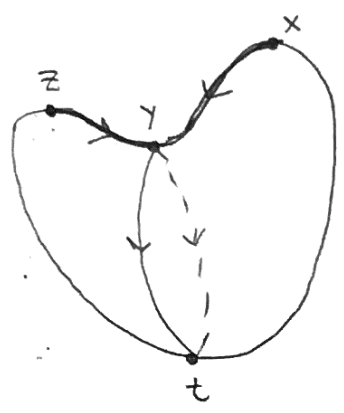
Q1: Is $\# \mathcal{M}(p, q)$ finite?

Q2: $(\partial^{MS})^2 = 0$?

A: $\mathcal{M}(p, q)$ can be compactified (for any p, q) by adding in "broken flow lines" \Rightarrow Q1 and Q2 are true

\leftarrow Sum of ∂ points of a compact 1-man = 0.

Example



$$C_2 \cong \mathbb{Z} \langle x, z \rangle \quad \partial x = \partial z = y$$

$$C_1 \cong \mathbb{Z} \langle y \rangle \quad \partial y = (+1 - 1) \cdot t = 0$$

$$C_0 \cong \mathbb{Z} \langle t \rangle$$

Remark: Unstable manifolds provide cell decomp. as described in beginning.

1982: Ed Witten: The Morse-Smale complex arises naturally from supersymmetric quantum mechanics on M .

Crash Course in QM: Particle on (M, g)

Classical Mech

- Phase Space T^*M
- $H: T^*M \rightarrow \mathbb{R}$
- Ham. Flow $\mathcal{F}: \mathbb{R} \rightarrow \text{Diff}(M)$

Quantization \rightsquigarrow

Quantum Mech

- Hilbert Space $\mathcal{H} = L^2(M)$
- $\hat{H}: \mathcal{H} \rightarrow \mathcal{H}$ self-adjoint
- $\mathbb{R} \rightarrow \text{U}(\mathcal{H})$
- $t \mapsto e^{iHt/\hbar}$

QM

Rep. of $\mathbb{R}\langle \hat{H} \rangle$

Susy QM

Rep. of $\mathbb{R}\langle \hat{H}, \hat{Q} \rangle$

$$\hat{H} = \hat{Q}^2$$

$$[\hat{H}, \hat{Q}] = 0$$

"2nd order diff. op. written as square of a 1st order op."

Ex: Free Particle on (M, g) (compact)

1) QM: $\mathcal{H} = \Omega^0(M)$; $\langle f, g \rangle = \int f \wedge *g$ Hodge star op: $\Omega^k \rightarrow \Omega^{n-k}$
 $\hat{H} = \Delta$ Hodge Laplacian
 $(= -g^{ij} \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j})$

2) Susy QM: $\mathcal{H} = \Omega^*(M) \otimes_{\mathbb{R}} \mathbb{C}$
 $\hat{H} = \Delta = (d + d^*)^2 : \Omega^k \rightarrow \Omega^k$
 $\hat{Q} = d + d^* : \Omega^k \rightarrow \Omega^{k+1} \oplus \Omega^{k-1}$

Natural Question: What is the Spectrum of \hat{H} ? ← "Stationary" States

Note: $\hat{H} = \hat{Q}^2 \Rightarrow$ Smallest eigenvalue ≥ 0

$$\text{Ker}(\hat{H}) = \{ \omega \in \Omega^* : \Delta \omega = 0 \} = \mathcal{H}^*(M)$$

harmonic forms on M

Hodge Theorem: $\mathcal{H}^*(M) \cong H_{dR}^*(M)$

(as \mathbb{Z} -graded vector spaces).

Now Consider an $\mathbb{R}_{\geq 0}$ -Family of deformations of \hat{H} : "Add a potential"

$$\hat{H}_s = (d_s + d_s^*)^2$$

$$d_s := d + s dh \wedge (-) : \Omega^k \rightarrow \Omega^{k+1}$$

where

- $h : M \rightarrow \mathbb{R}$ Morse
- $s \in \mathbb{R}_{>0}$

Q: How does $\text{Ker}(\hat{H}_s)$ change as I vary s ?

Observation:

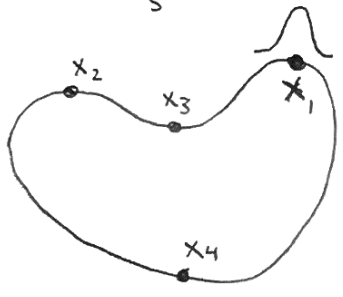
$$d_s = e^{sh} \circ d \circ e^{-sh} \leftarrow \text{mult. by } e^{-sh}$$

\Rightarrow The map $\Omega^k \rightarrow \Omega^{k+1}$
 $v \mapsto e^{sh} v$

- 1) Gives an iso. of complexes $(\Omega^k, d_s) \xrightarrow{\sim} (\Omega^k, d)$
 - 2) " " " $\text{Ker}(\hat{H}_s) \xrightarrow{\sim} \text{Ker}(\hat{H})$
- $\Rightarrow \text{Ker}(\hat{H}_s) \cong H^1_{\text{dR}}(M)$

The $s \rightarrow \infty$ limit:

Guess at v s.t. $\hat{H}_s v = 0$:



$$\left(\begin{aligned} \hat{H}_s v &= 0 \\ \Rightarrow d_s v &= 0 \\ \Rightarrow v &\text{ localized around } \text{Crit}(h) \end{aligned} \right)$$

Claim: $d_s x_i \neq 0$!

Why? "Instanton Corrections"

$\{l_1, \dots, l_r, \psi_1, \dots\}$
 orthonormal basis

$$d_s x_i = \sum_{l \in \text{Crit}(h)} \langle l, d_s x_i \rangle l + \underbrace{\sum_{k=1}^{\infty} \langle \psi_k, d_s x_i \rangle \psi_k}_C$$

6

Claim: • $\langle \ell, d_S X_i \rangle = \text{Signed Sum over gradient Flow lines } X_i \rightarrow \ell + \mathcal{O}(\frac{1}{S})$

• $C = \mathcal{O}(\frac{1}{S})$

Thus,

$$(\Omega^{n-0}, d_S) \xrightarrow{S \rightarrow \infty} (C^{\text{MS}} \otimes_{\mathbb{Z}} \mathbb{C}, \partial^{\text{MS}})$$

and we expect

$$H_{\text{DR}}^{n-0}(M) \otimes_{\mathbb{R}} \mathbb{C} \cong H_{\bullet}^{\text{MS}}(M) \otimes_{\mathbb{Z}} \mathbb{C} \cong H_{\bullet}^{\text{sing}}(M; \mathbb{C})$$

What are instanton corrections?


Quantum Mechanics on $M \cong \mathbb{R}$ w/ potential $U: M \rightarrow \mathbb{R}$

Quantum Hamiltonian:

$$\hat{H} = -\frac{1}{\hbar^2} \frac{\partial^2}{\partial x^2} + U(x)$$

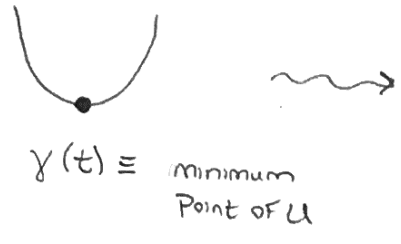
What are its lowest eigenstates?

(See 5 old)

Ex: $U =$ 

Lowest energy Eigenstates : $\hat{H}\psi = E_{low}\psi$

Classically:



Quantum Mechanically 



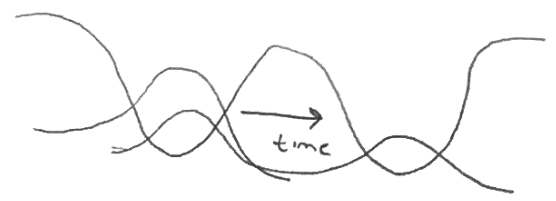
Guess:



and



But in reality: tunneling!



$\psi_1 + \psi_2$ (or $\psi_1 - \psi_2$) is the actual ground state.

Path Integral Point of View: Why is $e^{-iHt}\psi_1 \neq \psi_1$?

Answer: Quantum Corrections From "Instantons"